

Transverse spin asymmetries in J/ψ production at COMPASS

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Faculty of mathematics
and physics





- 1 Nucleon structure
- 2 COMPASS experiment
- 3 TSAs in J/ψ lepton production
- 4 TSAs in $\pi^- p^\uparrow$ scattering
- 5 Conclusions



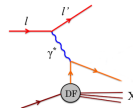
- **Parton model**

- Born in the late 60's,
- to describe electron-proton scattering $e(k) + p(P) \rightarrow e(k') + X$

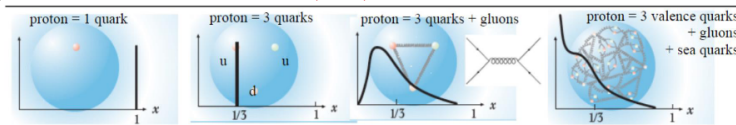
$$Q^2 = -q^2 = -(k' - k)^2, \quad x = \frac{Q^2}{2P \cdot q}$$

- Deeply-Inelastic Scattering (DIS) limit and “infinite momentum frame”:
 $P \rightarrow \infty, Q^2 \rightarrow \infty, x$ stays finite.

$$\frac{d\sigma}{dx dQ^2} = \frac{4\pi\alpha^2}{Q^4} D(x, Q^2) \sum_i e_i^2 f_1^i(x)$$



- **Point-like constituents (partons)** with momentum and charge: $k = xP, e = 2/3, -1/3$.
- $f_1(x)$ – **Parton Distribution Function (PDF)**,



- Drell-Yan reaction (1970)

$$\frac{d\sigma}{dx_a dx_b} = \frac{4\pi\alpha^2}{9q^2} \sum_i e_i^2 f_1^i(x_a) f_1^{\bar{i}}(x_b)$$

- **QCD**

- $f = f(x, Q^2)$, but the dependence on Q^2 is calculable (DGLAP equations).



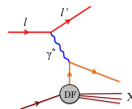
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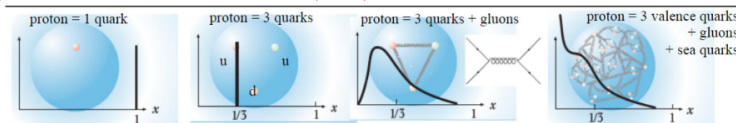
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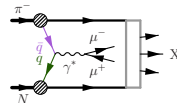


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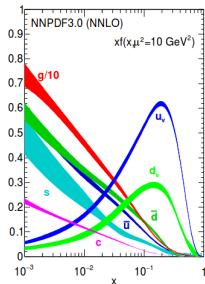
Number density



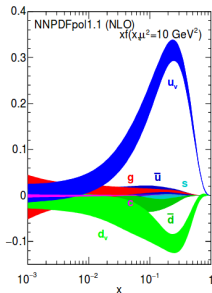
Helicity



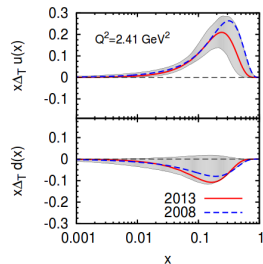
Transversity



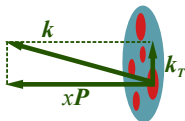
NNPDF 3 number density



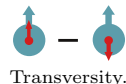
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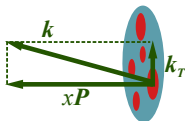
Transversity from Anselmino *et al.*



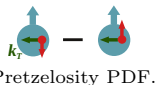
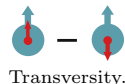
- If parton intrinsic k_T is not integrated over,
- “three-dimensional” objects $f(x, k_T^2, Q^2)$.
- Accessible in
 - semi-inclusive deep-inelastic scattering (SIDIS),
 - Drell–Yan dilepton production,
 - proton–proton collisions...



		Parent hadron polarization		
		Unpolarised	Longitudinal	Transverse
Parton polarisation	U	$f_1(x, k_T^2)$ (number density)		$f_{1T}^\perp(x, k_T^2)$ (Sivers)
	L		$g_1(x, k_T^2)$ (helicity)	$g_{1T}(x, k_T^2)$ (Kotzinian–Mulders)
	T	$h_1^\perp(x, k_T^2)$ (Boer–Mulders)	$h_{1L}^\perp(x, k_T^2)$ (worm-gear)	$h_1(x, k_T^2)$ (transversity) $h_{1T}^\perp(x, k_T^2)$ (pretzelocity)



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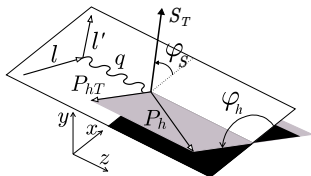


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- Semi-Inclusive Deep-Inelastic Scattering (SIDIS) off polarised nucleons

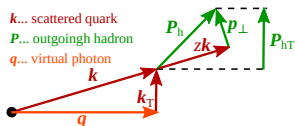
$$\ell + \vec{N} \rightarrow \ell' + h + X, \quad \ell + N^\uparrow \rightarrow \ell' + h + X$$

- Q^2 , x , z (fraction of available energy transferred to h), \mathbf{P}_T (transverse momentum of h).



γN frame.

- \mathbf{k}_T is not directly observable, only convolutions of TMD PDFs and fragmentation functions (FFs) over \mathbf{k}_T and \mathbf{p}_\perp .
- Measured since ≈ 2000 : HERMES, COMPASS, JLab.



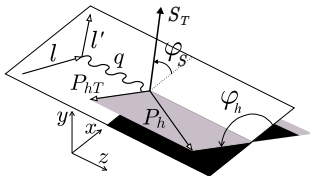
Transverse momenta
in target rest frame.



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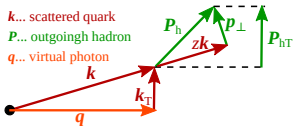
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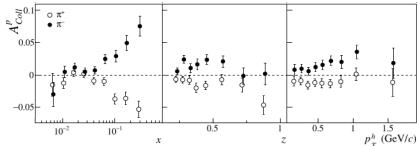


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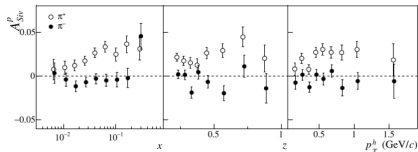


Transverse momenta in target rest frame.



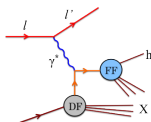
Collins asymmetry at COMPASS.

Arises from transverse polarisation of q in p^\uparrow (transversity PDF h_1), and anisotropic fragmentation of q^\uparrow (Collins TMD FF H_1^\perp).



Sivers asymmetry at COMPASS.

From correlation of \mathbf{k}_T of q with the spin of p^\uparrow (Sivers TMD PDF f_{1T}^\perp)



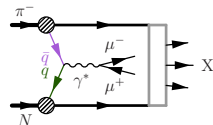
SIDIS on transversely polarised nucleons

- Structure functions F :

$$F = \text{PDF}_{q,p} \otimes \text{FF}_{q \rightarrow h}$$

- For example:

- $F_{UU}^{\cos \phi_h}$ and $F_{UU}^{\cos 2\phi_h}$ linked to $h_{1,P}^\perp$,
- $F_{UT,T}^{\sin(\phi_h - \phi_S)} = f_{1T,P}^\perp \otimes D_1$.
- $F_{UT}^{\sin(\phi_h + \phi_S)} = h_{1,P} \otimes H_1^\perp$,



Drell–Yan on transversely polarised nucleons

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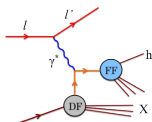
A sign change predicted for Sivers and Boer–Mulders functions:

$$f_{1T}^{\perp q} |_{\text{SIDIS}} = -f_{1T}^{\perp q} |_{\text{DY}}$$

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[J. Collins, Phys.Lett. B536

(2002) 43]



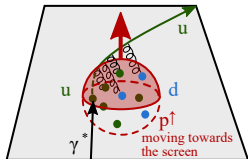
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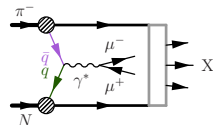
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Siverts effect in SIDIS

(as described by [M. Burkardt, Nucl.Phys. A735 (2004) 185].



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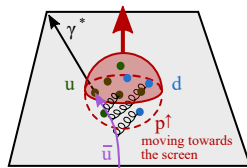
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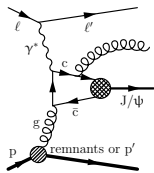
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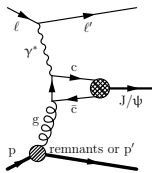


Siverts effect in Drell–Yan drawn in the same manner.

In lepton–nucleon scattering



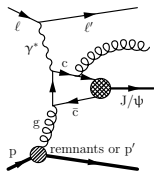
Color singlet model.



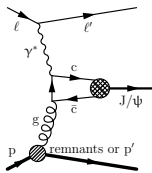
Color evaporation model (CEM) or color octet model (COM).

- **Color singlet model:** $c\bar{c}$ has to be in color singlet state to form the J/ψ .
- **CEM:** the color ‘evaporates’ from $c\bar{c}$ via soft gluon interactions.
- **COM:** NRQCD factorisation, different transition probabilities from different initial states (more free parameters than CEM).
- **These processes give access to gluon TMD PDFs via the ‘photon–gluon fusion’ (PGF).**
- **Diffractive production:** via exchange of a color-less particle.
- Also the diffractive production could be approached by perturbative QCD, but contains different information.
- **Feed-down:** decay of heavier charmonia – any information on nucleon structure is lost.

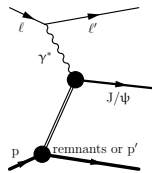
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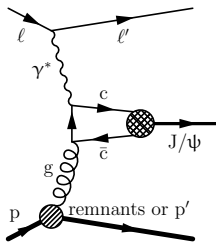


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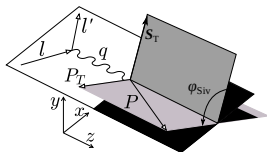
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In lepton-nucleon scattering



J/ψ formed in PGF.

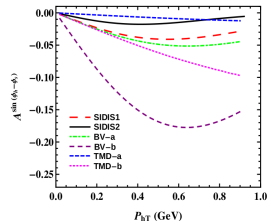
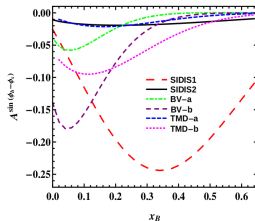


The Siverts angle
(here $P = P_{J/\psi}$).

- Assuming PGF process (plus CEM or COM),
- azimuthal distribution of g preserved, $\phi_g = \phi_{J/\psi}$,
- \rightarrow Siverts-like modulation in φ_{Siv}

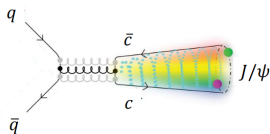
$$\sigma(\phi_{Siv}) = \sigma_0 (1 + f P_{tar.} A_{Siv}^P \sin(\varphi_{Siv}))$$

- [Mukherjee and Rajesh, Eur.Phys.J.C77 (2017)],
[Bacchetta *et al.*, Eur.Phys.J.C80 (2020)].

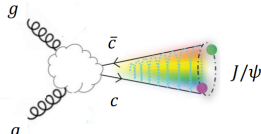


Projections for COMPASS energy,
[Mukherjee and Rajesh, Eur.Phys.J.C77 (2017)].

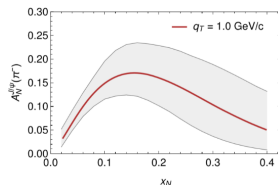
In pion–nucleon scattering



Quark–antiquark annihilation¹.



Gluon–gluon fusion¹,
the cloud = several perturbative contributions.



[Anselmino *et al.* Phys.Lett.B770 (2017)]

- The rainbow area represents CEM or COM J/ψ production from $c\bar{c}$.
- $q\bar{q}$ annihilation: access to quark TMD PDFs.
 - A large Sivers asymmetry in $\pi^- p^\uparrow \rightarrow \mu^- \mu^+ X$ at COMPASS was predicted, assuming only $q\bar{q}$ [Anselmino *et al.* Phys.Lett.B770 (2017)].
- gg fusion: access to gluon TMD PDFs.
 - For example, to d-type Sivers function if produced in $\pi^- p^\uparrow$ scattering.
- Feed-down: decay of heavier charmonia – any information on nucleon structure is lost.
- The result is a mix of the processes. The ratio depends on J/ψ production mechanism.
 - Studies suggest that gg fusion dominates at COMPASS [Chang *et al.*, Phys.Rev.D102 (2020)]

¹Diagrams: courtesy of Pietro Faccioli.



[S. Arnold *et al.*, Phys.Rev.D79 (2009) 034005]

$$\frac{d\sigma}{d\Omega} \propto (F_U^1 + F_U^2) (1 + A_U^1 \cos^2 \theta_{CS})$$

$$\times \left\{ \begin{aligned} & 1 + D_{[\sin^2 \theta_{CS}]} A_U^{\cos 2\varphi_{CS}} \cos 2\varphi_{CS} + D_{[\sin 2\theta_{CS}]} A_U^{\cos \varphi_{CS}} \cos \varphi_{CS} \\ & + S_T \left[\begin{aligned} & A_T^{\sin \varphi_S} \sin \varphi_S \\ & + D_{[\sin 2\theta_{CS}]} \left(\begin{aligned} & A_T^{\sin(\varphi_{CS} - \varphi_S)} \sin(\varphi_{CS} - \varphi_S) \\ & + A_T^{\sin(\varphi_{CS} + \varphi_S)} \sin(\varphi_{CS} + \varphi_S) \end{aligned} \right) \\ & + D_{[\sin^2 \theta_{CS}]} \left(\begin{aligned} & A_T^{\sin(2\varphi_{CS} - \varphi_S)} \sin(2\varphi_{CS} - \varphi_S) \\ & + A_T^{\sin(2\varphi_{CS} + \varphi_S)} \sin(2\varphi_{CS} + \varphi_S) \end{aligned} \right) \end{aligned} \right\}$$

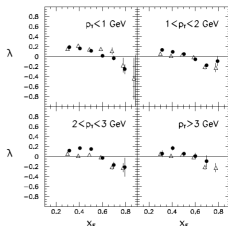
$$D_{[f(\theta_{CS})]} = f(\theta_{CS}) / (1 + A_U^1 \cos^2 \theta_{CS})$$

Cross-section with unpolarised target:

$$\frac{d\sigma}{d\Omega} = \frac{3}{4\pi} \frac{1}{\lambda + 1} \left(1 + \lambda \cos^2 \theta_{CS} + \mu \sin 2\theta \cos \varphi_{CS} + \frac{\nu}{2} \sin^2 \theta \cos 2\varphi_{CS} \right)$$

$$\lambda = A_U^1, \mu = A_U^{\cos \varphi_{CS}} \text{ and } \nu = 2A_U^{\cos 2\varphi_{CS}}$$

- The same parametrisation describes Drell–Yan and J/ψ cross-section.
- The interpretation of the structure functions F differs.
- ‘Naive’ Drell–Yan model: $\lambda = 1, \mu = \nu = 0$.
- Lam–Tung relation for Drell–Yan: $\lambda + 2\nu = 1$.
- λ plays role in the kinematic factors D .
- Drell–Yan: $\lambda \approx 1$ from experiments.
- J/ψ : kinematically dependent.

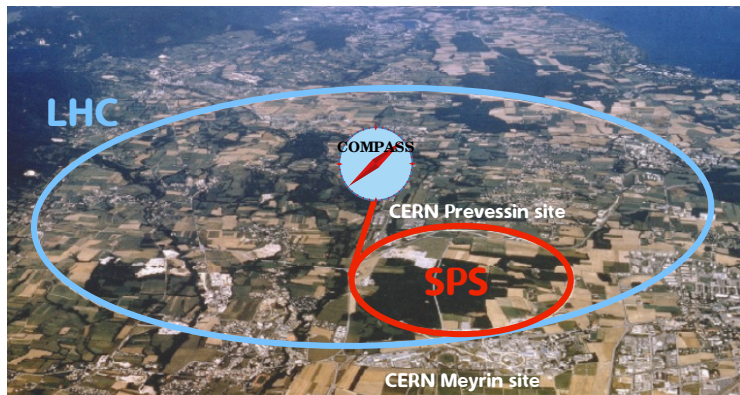


λ measured by NuSea experiment

[NuSea, Phys.Rev.Lett. 91 (2003)].



- COMPASS Collaboration: 24 institutions from 13 countries (≈ 220 physicists).
- Experimental area: CERN Super Proton Synchrotron (SPS) North Area.
- Multi-purpose apparatus with rich physics program since 2002 aimed at hadron structure and spectroscopy.





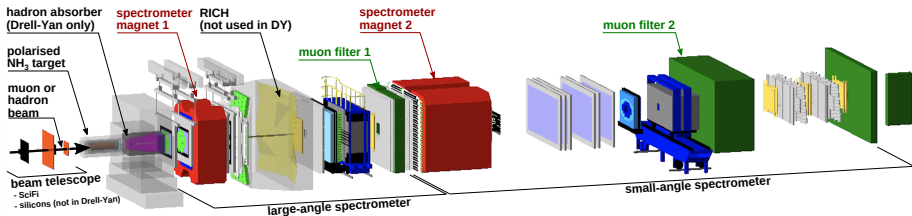
- Large polarised solid-state target with 2 or 3 oppositely-polarised cells.
- Two-stage spectrometer, about 350 detector planes, μ identification.

SIDIS with transversely-polarised target

- 2002–2004 with d^\uparrow (${}^6\text{LiD}$, old magnet).
- 2007 and 2010 with p^\uparrow (NH_3 , new mag.)
- 2021–2022 with d^\uparrow (${}^6\text{LiD}$, new mag.)
- 160 GeV/c μ^+ beam
(about $3.5 \times 10^8 \mu/\text{spill}$ of 10 s).
- Triggering on the scattered μ .

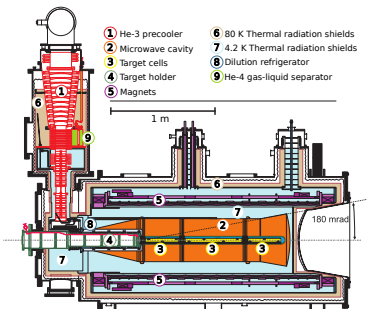
Drell–Yan with transversely-polarised target

- 2015 and 2018 with p^\uparrow (NH_3 , new mag.)
- 190 GeV/c π^- beam
(about $10^9 \pi/\text{spill}$ of 10 s).
- With a hadron absorber.
- Triggering on 2μ .



COMPASS Drell–Yan setup.

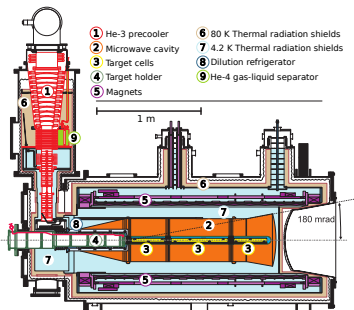
The SIDIS setup has beam momentum stations, silicon beam telescope and no hadron absorber.



Polarised target (new magnet).

Large solid-state polarised target

- Super-conducting magnets:
 - 2.5 T solenoid, 0.6 T dipole.
 - Old (SMC) target magnet: 70 mrad acceptance.
 - New (COMPASS) target magnet: 180 mrad.
- MW system for dynamic nuclear polarisation.
- Polarisation is measured by NMR.
- Dilution refrigerator → frozen spin mode at 70 mK.
- The target contains also unpolarised nuclei
 - Dilution of the signal by fP_N .
 - f : fraction of cross-section on polarisable nuclei.
 - P_N : polarisation of the polarisable nuclei.
 - Polarised d: ${}^6\text{LiD}$, $f = 0.4$, $P_d = 0.5$
 - Polarised p: NH_3 , $f = 0.16$, $P_p = 0.9$
- Acceptance in polarisation-dependent azimuthal angles is cancelled in combinations of target cells and data taking periods.



Polarised target (new magnet).



Operation with 2 cells.



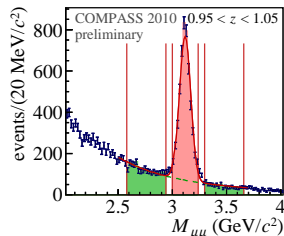
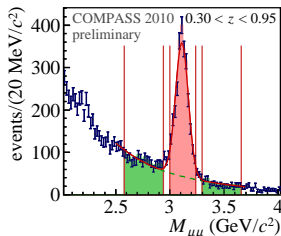
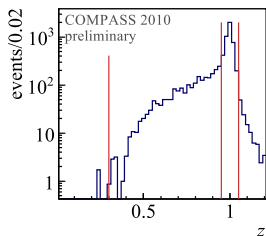
Operation with 3 cells.

Large solid-state polarised target

- Super-conducting magnets:
 - 2.5 T solenoid, 0.6 T dipole.
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 - New (COMPASS) target magnet: 180 mrad.
- MW system for dynamic nuclear polarisation.
- Polarisation is measured by NMR.
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 - Polarised p: NH_3 , $f = 0.16$, $P_P = 0.9$
- **Acceptance in polarisation-dependent azimuthal angles is cancelled** in combinations of target cells and data taking periods.



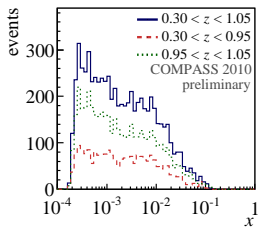
- $\mu^+p^\uparrow \rightarrow \mu^+J/\psi X \rightarrow \mu^+\mu^+\mu^-X$.
- Both possible combinations of $\mu^+\mu^-$ used.
- 2010 proton data.
- No Q^2 cut imposed (hard scale = $c\bar{c}$ mass).
- Two bins z : inclusive, **exclusive**.
- Clear J/ψ signal (3.1 GeV/c^2 , $\sigma \approx 55 \text{ MeV}/c^2$).
- Small background, limited statistics (≈ 2300 incl., 4500 excl.).



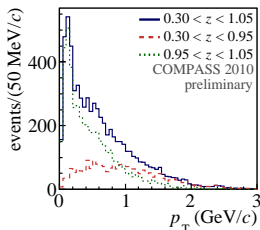
Energy fraction transferred to J/ψ .

Invariant mass, inclusive bin.

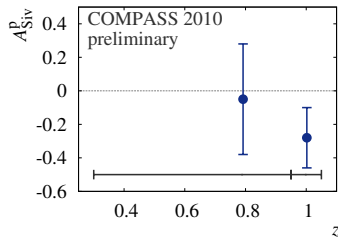
Invariant mass, exclusive bin.



The x_{Bj} distribution.

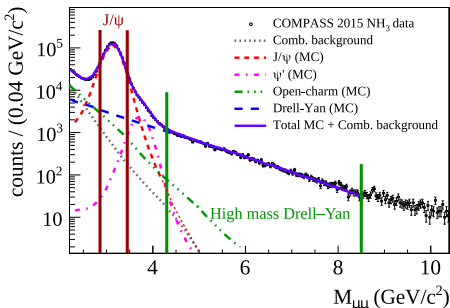


The P_T distribution.

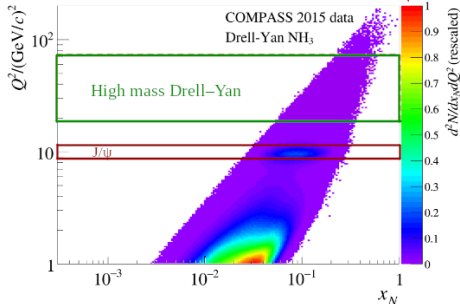


The measured Siverson-like asymmetry.

- $A_{Siv}^P = -0.28 \pm 0.18$ (**preliminary**, exclusive J/ψ).
- Prospects for improving statistics:
 - e^+e^- channel: spectrometer not optimal for electrons, probably high background...
 - 2002–2004 ^6LiD data: rather small statistics,
 - 2007 NH_3 data could bring something,
 - Planned 2022 ^6LiD data: ≈ 2010 statistics.
- We are considering analysing other J/ψ asymmetries and writing a paper.

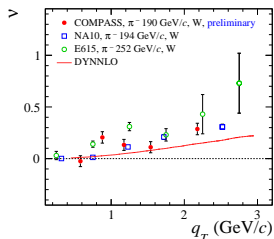
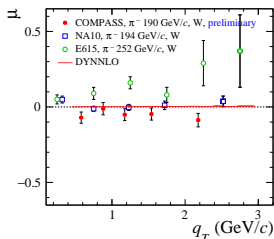
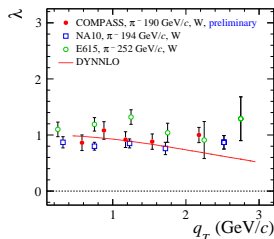


2015 data and reconstructed MC.

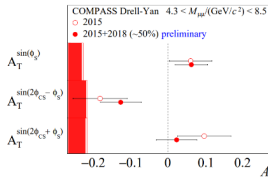
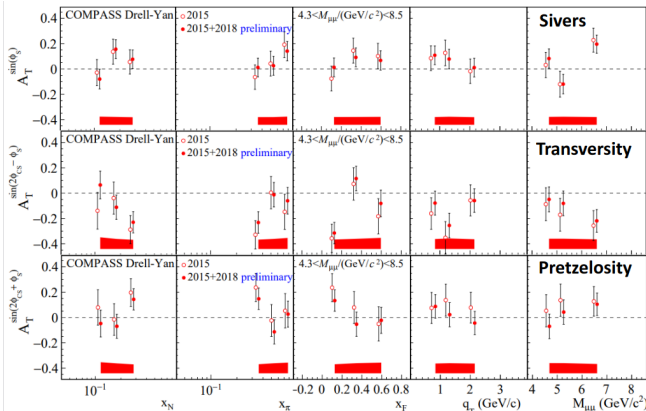


Kinematic coverage in x_N and Q^2

- $\pi^- p^\uparrow \rightarrow \mu^+ \mu^-$.
- Invariant mass distribution is smeared by the hadron absorber.
- Combinatorial background evaluated from like-sign $\mu\mu$ in the data.
- Open-charm background evaluated from Monte Carlo.
- $M_{\mu\mu} \in [4.3, 8.5]$ GeV/ c^2 : High mass Drell-Yan region (96% pure Drell-Yan)
 - TSAs from 2015 data published [COMPASS, Phys.Rev.Lett.119(11), 112002 (2017)].
- $M_{\mu\mu}$ in J/ψ region: more than 90% pure J/ψ , depending on the precise cut.
 - Ongoing analysis of the TSAs.
 - About $30\times$ more data with respect to high-mass Drell-Yan.



- Unpolarised asymmetries in line with previous experiments.
- In line with Lam–Tung relation, within uncertainties.
- Obtained using 2018 data (better Monte Carlo description than 2015).
- J/ψ analysis is ongoing.



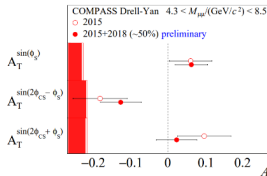
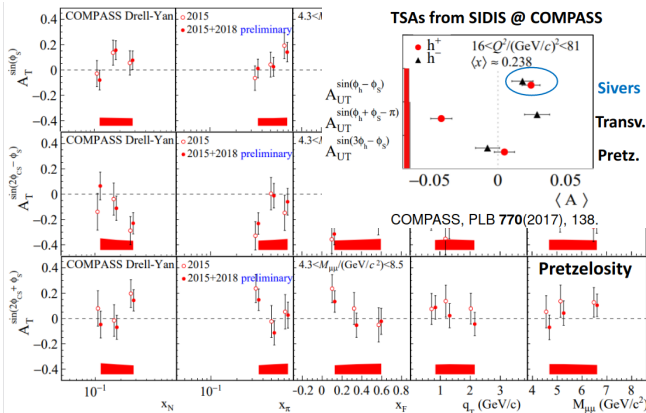
Sivers $\sim 1\sigma$ above zero

Transversity $\sim 2\sigma$ below zero

Pretzelocity $\sim 1\sigma$ above zero

TSAs in the high-mass Drell–Yan range [COMPASS, Phys.Rev.Lett.119(11), 112002 (2017)].

- The results support the sign-change prediction, although with a limited precision.
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- It can thus address the universality of the TMDs.
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Thank you for your attention!